

* wave particle dualism- As light dual nature in certain phenomena light behave like a wave. Such as interference deflection in certain other phenomena it behave like a particle. such as photo electric effect and Compton effect.

* De- Broglie introduce that matter also should have dual nature. The wave length of that wave is given by

$$\lambda = \frac{h}{p} = \frac{h}{mv}$$

where v = velocity of particle
 h = plank's constant = 6.63×10^{-34}
 m = mass of particle

Imp De-Broglie Hypothesis- was Based upon the following effects.

- (i) Light and a matter both are the form of energy and they can be interchange.
- (ii) As Light has dual nature so, matter also should have dual nature.

"According to De-Broglie Hypothesis - a wave is associated with moving particle".

The wave length of the wave associated with moving particle of mass (m) moving with velocity (v) as given by-

$$\lambda = \frac{h}{p} = \frac{h}{mv}$$

Expression for De-Broglie wave length-

According to Einstein's mass energy relation the energy is defined as.

$$E = mc^2 \quad \text{--- (1)}$$

According to Energy of Photon is also defined as

$$E = h\nu$$

$$E = h \frac{c}{\lambda}$$

--- (2)

$$\{c = \nu \lambda\}$$

From eq (1) and eq (2) we have

$$h\nu = mc^2$$

$$m = \frac{h}{\lambda}$$

$$\lambda = \frac{h}{m}$$

If the moving particle is material particle then

$$\lambda = \frac{h}{mv}$$

$$\lambda = \frac{h}{p}$$

$$\{p = mv\}$$

c-918

Different Expression for De-Broglie wave length-

Wavelength

The K.E. of moving particle of mass (m) moving with velocity (v) is given by

$$E = \frac{1}{2} mv^2$$

$$\{K.E. = E\}$$

$$E = \frac{(mv)^2}{2m}$$

$$E = \frac{p^2}{2m}$$

$$p = \sqrt{2mE}$$

(2)

If the moving particle is a charged particle of charge (q) accelerated by a potential difference then

$$E = qV$$

$$\lambda = \frac{h}{\sqrt{2mqV}}$$

(3)

If the moving particle as n electron-

$$\lambda = \frac{h}{\sqrt{2mnev}}$$

$$\lambda = \frac{6.63 \times 10^{-34}}{\sqrt{2 \times 9.1 \times 10^{-31} \times 1.6 \times 10^{-19} \times v}}$$

$$\lambda = \frac{12.28}{\sqrt{V}} \text{ \AA}$$

If the moving particle is Neutron then

$$\lambda = \frac{h}{\sqrt{2m_n E}}$$

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$$\lambda = \frac{h}{\sqrt{2m_n E}}$$

$$\lambda = \frac{6.63 \times 10^{-34}}{\sqrt{2 \times 1.67 \times 10^{-27} \times 1.6 \times 10^{-19} \times E}}$$

$$\lambda = \frac{6.63 \times 10^{-34}}{\sqrt{2 \times 1.67 \times 10^{-27} \times 1.6 \times 10^{-19} \times E}}$$

$$\lambda_n = \frac{0.286 \text{ \AA}}{\sqrt{E_{\text{eV}}}}$$

Q) If the particle is in thermal equilibrium then

$$E = \frac{3}{2} kT$$

$$\lambda = \frac{h}{\sqrt{2m \times \frac{3}{2} kT}}$$

$$\lambda = \frac{h}{\sqrt{3mkT}}$$

Que. $v = 100 \text{ V}$ $^2\text{He}^4$

$$\lambda = \frac{h}{mv}$$

$$\lambda \lambda = \frac{h^2}{2mqv}$$

$$\lambda \lambda = \frac{6.63 \times 10^{-34}}{2 \times 1.6 \times 10^{-19} \times 100}$$

$$\lambda \lambda = 0.0017 \text{ \AA}$$

* properties of de-Broglie matter wave - De-Broglie matter wave have the following characteristics -

→ larger the velocity of the particle smaller the wave length of the wave associated with it and vice-versa.

→ when $\theta = 0$ then $\lambda = \infty$ that is the wave is associated only with moving particle.

→ smaller the mass of particle, larger the wavelength of the wave associated with it and vice-versa.

→ Both aspects never exist together in same experiment.

→ matter wave are not electromagnetic wave.

Expression for the velocity of De-Broglie matter wave - The Energy of material particle can be expressed as -

$$E = mc^2$$

The energy is also defined as - $E = h\nu$

$$E = mc^2 = h\nu$$

$$\nu = \frac{mc^2}{h} \quad \text{--- (1)}$$

The wavelength of De-Broglie matter wave is given by

$$\lambda = \frac{h}{mv} \quad \text{--- (2)}$$

where, ν = velocity of particle

velocity of De-Broglie matter wave or wave velocity

(Phase velocity) is defined as

$$\nu_p = \nu \lambda$$

From eq (1) and eq (2)

$$\nu_p = \frac{mc^2}{h} \times \frac{h}{mv}$$

$$\nu_p = \frac{c^2}{v}$$

From here it is clear that the phase velocity is greater than the speed of light.

Imp (ii) wave velocity (v_p) - Phase velocity is defined as the ratio of angular velocity to the per propagation constant.

$$v_p = \frac{\omega}{k}$$

(iii) equation of wave can be written as -

$$Y = A \sin(\omega t - kx)$$

$$\therefore \omega t - kx = \text{const.}$$

diff w.r. to t we have

$$\omega - k \frac{dx}{dt} = 0$$

$$\frac{dx}{dt} = \frac{\omega}{k}$$

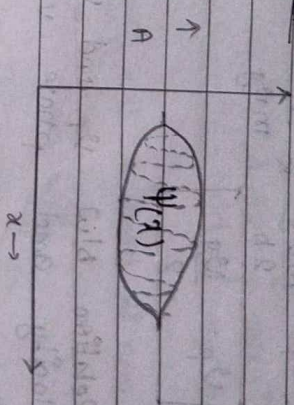
That is the velocity of constant phase of a wave is known as phase velocity.

Group Velocity - when a group of wave propagates along the same direction then the observed velocity is the velocity of maximum amplitude which are known as Group velocity.

That is
$$v_g = \frac{d\omega}{dk}$$

True, the group velocity is define as the ratio of change in angular velocity to the change in propagation constant.

Wave packet -



When a group of wave propagate along the same direction having amplitude and phase search that they interfere constructively over a small region of space. where the particle can be located is known as wave packet.

The velocity of wave packet is also known as group velocity.

Relation b/w v_p and v_g for Non-relativistic free particle

Energy of Non-relativistic free particle can be express as.

$$E = h\nu = \frac{1}{2}mv^2$$

$$\nu = \frac{mv^2}{2h} \quad \text{--- (1)}$$

The wavelength is define as -

$$\lambda = \frac{h}{mv} \quad \text{--- (2)}$$

The wave velocity is given by -

$$v_p = \frac{v}{\lambda}$$

then

$$v_p = \frac{m v g}{2h} \times \frac{h}{m v g}$$

$$v_p = \frac{v g}{2}$$

Imp # General Relation b/w v_p and $v_g \rightarrow$ The group velocity and phase velocity are define as-

$$v_g = \frac{d\omega}{dk} \quad \text{--- (1)}$$

$$v_p = \frac{\omega}{k} \quad \text{--- (2)}$$

from eqⁿ (1)

$$v_g = \frac{d\omega}{d(2\pi)} \quad \left\{ k = \frac{2\pi}{\lambda} \right\}$$

$$v_g = \frac{d\omega}{d\lambda}$$

$$v_g = \frac{2\pi(-\lambda^{-2} d\lambda)}{2\pi \frac{d\omega}{d\lambda}} \quad \text{--- (3)}$$

from eqⁿ (2) we have -

$$\omega = k \cdot v_p$$

$$\omega = \frac{2\pi v_p}{\lambda}$$

diff w.r. to λ

Case-I If the medium is ~~known~~ ϵ non-dispersive.

$$\frac{d v_p}{d\lambda} = 0$$

$$v_p = v_g$$

Imp

Que. show that

(i) $v_g = v$

(ii) $v_p v_g = c^2$

Solⁿ - proof (i)

The group velocity can be expressed as -

$$v_g = \frac{d\omega}{dk}$$

$$v_g = \frac{d\omega/dt}{dk/dt} \quad \text{--- (1)}$$

The angular velocity is define as

$$\omega = 2\pi f = \frac{2\pi E}{h} = \frac{2\pi(m c^2)}{h}$$

(2)

$$\omega = \frac{2\pi m c^2}{h} \left(1 - \frac{v^2}{c^2} \right)^{1/2}$$

$$\left\{ m = \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}}} \right\}$$

$$k = \frac{2\pi}{\lambda} = \frac{2\pi}{h} = \frac{2\pi m v}{h}$$

$$K = \frac{2\pi m_0 v}{h \left(1 - \frac{v^2}{c^2}\right)^{1/2}} \quad (3)$$

The diff eqn w.r. to v

$$\frac{dK}{dv} = \frac{2\pi m_0 c^2}{h} \frac{d}{dv} \left(1 - \frac{v^2}{c^2}\right)^{-1/2}$$

$$\frac{dK}{dv} = \frac{2\pi m_0 c^2}{h} \left[\frac{-1}{2} \left[1 - \frac{v^2}{c^2}\right]^{-3/2} \left[-\frac{2v}{c^2}\right] \right] \quad (4)$$

$$\frac{dK}{dv} = \frac{2\pi m_0 v}{h \left(1 - \frac{v^2}{c^2}\right)^{3/2}} \quad (4)$$

diff eqn (3) w.r. to v

$$\frac{dK}{dv} = \frac{2\pi m_0}{h} \frac{d}{dv} \left[\frac{v}{c^2} \left(1 - \frac{v^2}{c^2}\right)^{1/2} \right]$$

$$\frac{dK}{dv} = \frac{2\pi m_0}{h} \left[\frac{1}{c^2} \left(1 - \frac{v^2}{c^2}\right)^{1/2} - v \cdot \frac{1}{2} \left(1 - \frac{v^2}{c^2}\right)^{-1/2} \left[-\frac{2v}{c^2}\right] \right]$$

$$\frac{dK}{dv} = \frac{2\pi m_0}{h} \left[\frac{1}{c^2} \left(1 - \frac{v^2}{c^2}\right)^{1/2} + \frac{v^2}{c^2} \left(1 - \frac{v^2}{c^2}\right)^{-1/2} \right]$$

$$\frac{dK}{dv} = \frac{2\pi m_0}{h} \left[\frac{1 - \frac{v^2}{c^2}}{c^2} + \frac{v^2}{c^2} \left(1 - \frac{v^2}{c^2}\right)^{-1/2} \right]$$

$$\frac{dK}{dv} = \frac{2\pi m_0}{h} \left[\frac{1 - \frac{v^2}{c^2}}{c^2} + \frac{v^2}{c^2} \left(1 - \frac{v^2}{c^2}\right)^{-3/2} \right]$$

$$\frac{dK}{dv} = \frac{2\pi m_0 v}{h \left(1 - \frac{v^2}{c^2}\right)^{3/2}} \quad (5)$$

From eqn (4), (5) we get

$$v g = \frac{2\pi m_0 v}{h \left(1 - \frac{v^2}{c^2}\right)^{3/2}} \times h \left(1 - \frac{v^2}{c^2}\right)^{3/2} = 2\pi m_0 v$$

$$[v g = v]$$

Hence proved

proof (ii)

We know that -

$$v_p = \frac{c^2}{v}$$

$$v_g = \frac{c^2}{v_g} \quad \left\{ v = v_g \right\}$$

$$[v_p v_g = c^2] \text{ Hence proved}$$

Physical significance of wave function - The wave function ψ as such has no physical significance but $|\psi|^2$ gives the probability density of finding the particle over the given region of space at any time.

for N -acceptable wave function it should be continuous, finite and single valued everywhere.

$$\frac{\partial^2 \psi}{\partial x^2} = \text{finite}, \quad \frac{\partial \psi}{\partial x} = \text{continuous}$$

In the particle exist in the universe, Normalization condition

$$\int_{-\infty}^{\infty} |\psi|^2 dv = 1$$

This is known as Normalization condition.
 -1 If the particle does not exist -

$$\int_{-\infty}^{\infty} |\psi|^2 dx = 0$$

-1 If the wave function is a complex function then $|\psi|^2 = \psi \psi^*$ where ψ^* = Complex conjugate of ψ

Schrödinger's time independent and dependent wave eqn - [Time independent] -

The eqn of motion of a wave in differential form can be expressed as -

$$\nabla^2 \psi - \frac{1}{\psi^2} \frac{\partial^2 \psi}{\partial t^2} = 0 \quad \text{--- (1)}$$

The soln of eqn (1) is given by -

$$\psi = \psi_0 e^{-i\omega t} \quad \text{--- (2)}$$

diff. eqn (2) twice with r. to t.

$$\left[\frac{\partial \psi}{\partial t} = -i\omega \psi_0 e^{-i\omega t} \right] \quad \text{--- (3)}$$

$$\frac{\partial^2 \psi}{\partial t^2} = -\omega_0 (-i\omega) i\omega e^{-i\omega t}$$

$$\frac{\partial^2 \psi}{\partial t^2} = -\omega^2 \psi_0 e^{-i\omega t}$$

$$\frac{\partial^2 \psi}{\partial t^2} = -\omega^2 \psi \quad \text{--- (4) put in (1) we get}$$

$$\nabla^2 \psi + \frac{\omega^2}{\psi^2} \psi = 0 \quad \text{--- (5)}$$

$$\omega = 2\pi \nu \Rightarrow \frac{2\pi \nu}{\lambda} \quad \text{if } \omega \rightarrow \nu \lambda$$

$$\omega = \frac{2\pi}{\lambda} \quad \left\{ d = \frac{h}{p} \right\}$$

$$\frac{\omega}{\lambda} = \frac{2\pi \nu}{\lambda} = \frac{p}{h}$$

$$\frac{\omega}{\lambda} = \frac{p}{h} \quad \left\{ \lambda = \frac{h}{p} \right\}$$

$$\frac{\omega^2}{\lambda^2} = \frac{p^2}{h^2}$$

Now,

$$E - V = \frac{p^2}{2m}$$

$$p^2 = 2m(E - V)$$

then

$$\frac{\omega^2}{\lambda^2} = \frac{2m}{h^2} (E - V) \quad \rightarrow \text{put in eqn (5) we get}$$

$$\nabla^2 \psi + \frac{2m}{h^2} (E - V) \psi = 0 \quad \text{--- (6)}$$

This is known as time independent wave eqn.

Time dependent wave eqn -

Schrödinger's Time dependent wave eqn - This eqn is obtained by elementing Eψ from independent wave eqn using eqn (3).

Now,

$$\frac{\partial \psi}{\partial t} = -i\omega \psi_0 e^{-i\omega t} \quad \text{eqn (3)}$$

$$\frac{\partial \psi}{\partial t} = -i(2\pi\nu) \psi$$

$$\frac{\partial \psi}{\partial t} = -i \left(\frac{2\pi E}{h} \right) \psi$$

$$\frac{\partial \psi}{\partial t} = \frac{-iE}{h} \psi$$

$$\frac{\partial \psi}{\partial t} = E \psi$$

$$\frac{i\hbar}{-i\hbar} \frac{\partial \psi}{\partial t} = E \psi$$

$$\left[i\hbar \frac{\partial \psi}{\partial t} = E \psi \right] \text{--- (7) put in eqn (6)}$$

from eqn (6)

$$\nabla^2 \psi + \frac{2m}{\hbar^2} (E - V) \psi = 0$$

$$\nabla^2 \psi + \frac{2m}{\hbar^2} (i\hbar \frac{\partial \psi}{\partial t} - V \psi) = 0$$

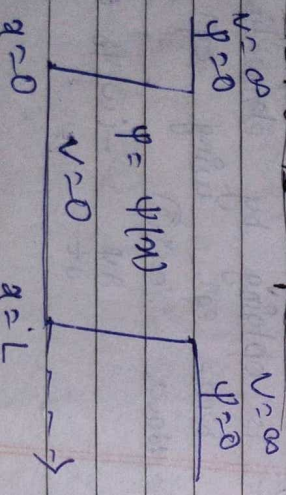
$$\left[i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + V \psi \right] \text{--- (8)}$$

$$[E\psi = H\psi]$$

$$H = \frac{-\hbar^2}{2m} \nabla^2 + V$$

Hamiltonian operator

Particle in one dimensional potential well (Box) -



When a particle moving along x-axis. The well then its potential energy is described by -

$$V = 0 \quad 0 \leq x \leq L$$

$$V = \infty \quad L \leq x \leq 0$$

The wave function for the particle is define as -

$$\psi = \psi(x) \quad 0 \leq x \leq L$$

$$\psi = 0 \quad L \leq x \leq 0$$

The motion of the particle inside the well is described by Schrodinger time independent wave equation -

$$\frac{\partial^2 \psi}{\partial x^2} + \frac{2m}{\hbar^2} (E - V) \psi = 0$$

$$\left[\frac{\partial^2 \psi}{\partial x^2} + \frac{2m}{\hbar^2} (E - V) \psi = 0 \right] \text{--- (1)}$$

The potential energy inside the well is zero

$$\left[\frac{\partial^2 \psi}{\partial x^2} + \frac{2m}{\hbar^2} E \psi = 0 \right] \text{--- (2)}$$

Let $\int \frac{2m}{\hbar^2} E \psi = k^2 \psi$ --- (3)

$$\left[\frac{\partial^2 \psi}{\partial x^2} + \psi k^2 = 0 \right] \text{--- (4)}$$

The solⁿ of eqn (4) is given by

$\psi = A \sin kx + b \cos kd$ -5

Now, using boundary conditions we have -

$x=0$ then $\psi=0$

put eq 1

$0 = A \sin 0 + b \cos 0$

$b + 0 = 0$

$b = 0$

Again using boundary condition we have -

$x=L$ then $\psi=0$ put eq 1

$0 = A \sin kL + b \sin kL$

$A \sin kL = 0$

$A \neq 0$ then $\sin kL = 0$

$\sin kL = \sin(\pm n\pi)$

$kL = \pm n\pi$

$k = \pm \frac{n\pi}{L}$ -6

$k^2 = \frac{n^2 \pi^2}{L^2}$ -7

put we have taken $k^2 = \frac{8\pi^2 m E}{h^2}$

$\frac{8\pi^2 m E}{h^2} = \frac{n^2 \pi^2}{L^2}$

$E_n = \frac{n^2 h^2}{8mL^2}$ -8

from here it is clear that the particle have the discrete energy value which are known as Eigen value. ? eigenvalue

then, $E_1 = \frac{h^2}{8mL^2}$

$E_2 = \frac{4h^2}{8mL^2} = 4E_1$

$E_3 = \frac{9h^2}{8mL^2} = 9E_1$

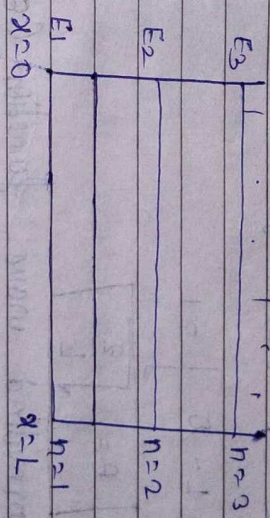


Fig- Eigen value diagram

The eigen function corresponding to each eigen value is called eigen function.

from eq 5 we have -

eq 1, 6 $\left\{ \begin{array}{l} k = \pm \frac{n\pi}{L} \\ \psi_n = A \sin kx \\ \psi_n = A \sin\left(\pm \frac{n\pi x}{L}\right) \end{array} \right.$ Eigen function
 -9

To normalize the wave function we used normalized condition.

Then $\int_0^L |u_n|^2 dx = 1$

$\int_0^L A^2 \sin^2 \left(\pm \frac{n\pi x}{L} \right) dx = 1$

$\left. \begin{aligned} \cos 2\theta &= 1 - 2\sin^2\theta \\ \sin^2\theta &= \frac{1 - \cos 2\theta}{2} \end{aligned} \right\}$

$A^2 \int_0^L \left[1 - \cos \left(\pm \frac{2n\pi x}{L} \right) \right] dx = 1$

$\frac{A^2}{2} \int_0^L \left[2 - \left(\frac{L}{\pm n\pi} \right) \sin \left(\pm \frac{2n\pi x}{L} \right) \right] dx = 1$

$\frac{A^2}{2} [L - 0] = 1$

$A = \sqrt{\frac{2}{L}}$

then normalized wave function (eigen function)

$u_n = \sqrt{\frac{2}{L}} \sin \left(\pm \frac{n\pi x}{L} \right)$

$u_0 = \sqrt{\frac{2}{L}} \sin \left(\pm \frac{0\pi x}{L} \right)$

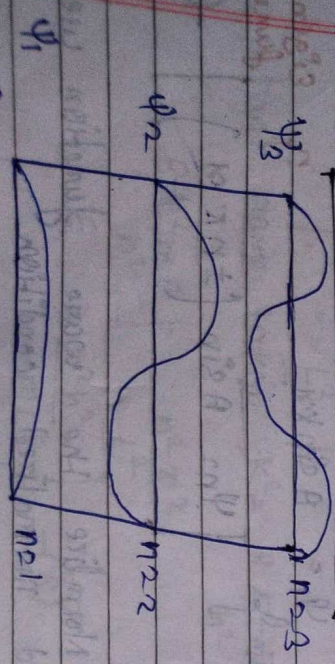


Fig- Eigen function Diagram

Compton effect-

When x-ray incident on a crystal then there are two types of radiation occurs in scattered region.

(1) which have larger wavelength is called secondary radiation and another radiation which have same wavelength as the wavelength of incident photon is called unmodified radiation.

The Compton shift depends only upon the angle of scattering.

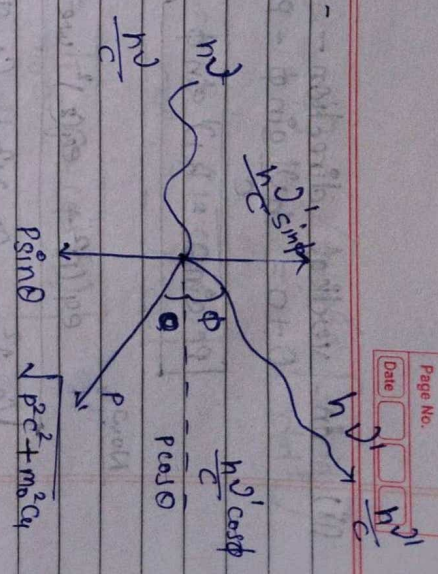
i.e. it does not depend upon the wavelength of incident photon.

(ii) In horizontal direction -

$E = \frac{h\nu}{c} = \frac{mc^2}{c} = mc$

$h\nu + 0 = \frac{h\nu'}{c} \cos \phi + p \cos \theta$

$p \cos \theta = h\nu - h\nu' \cos \phi$ (1)



(ii) In vertical direction -
 $0 + 0 = \frac{h\nu}{c} \sin \phi - p \sin \theta$

$$pc \sin \theta = h\nu \sin \phi \quad \text{--- (2)}$$

Now, eq (1) + eq (2) we get -

$$[pc]^2 = (h\nu)^2 + (h\nu)^2 - 2h\nu h\nu \cos \phi \quad \text{--- (3)}$$

Acc. to conservation law of energy -

$$h\nu + m_0 c^2 = h\nu' + [pc]^2 + m_0^2 c^4$$

$$[h\nu - h\nu'] + m_0 c^2 = [pc]^2 + m_0^2 c^4$$

$$[h\nu]^2 + (h\nu')^2 - 2h\nu h\nu' + m_0^2 c^4 + 2m_0 c^2 (h\nu - h\nu') = p^2 c^2 + m_0^2 c^4$$

$$p^2 c^2 = (h\nu)^2 + (h\nu')^2 - 2h\nu h\nu' + 2m_0 c^2 (h\nu - h\nu') \quad \text{--- (4)}$$

from eq (2) and (4) we get

$$(h\nu)^2 + (h\nu')^2 - 2h\nu h\nu' \cos \phi = (h\nu)^2 + (h\nu')^2 - 2h\nu h\nu' + 2m_0 c^2 (h\nu - h\nu')$$

$$h\nu h\nu' [1 - \cos \phi] = m_0 c^2 [h\nu - h\nu']$$

$$\frac{hc}{\lambda \lambda'} [1 - \cos \phi] = m_0 c^2 \left[\frac{c - c\lambda'}{\lambda} \right] \quad \left\{ \begin{array}{l} \nu = \frac{c}{\lambda} \\ \nu' = \frac{c}{\lambda'} \end{array} \right.$$

$$\frac{hc}{\lambda \lambda'} [1 - \cos \phi] = h m_0 c^2 \left[\frac{\lambda - \lambda'}{\lambda \lambda'} \right]$$

$$\lambda' - \lambda = \frac{h}{m_0 c} (1 - \cos \phi)$$

$$\Delta \lambda = \frac{h}{m_0 c} (1 - \cos \phi)$$

Compton shift.

Case-I when $\phi = 0$ then $\Delta \lambda = 0$ i.e. unmodified radiation is obtain in the direction of incident photon

Case-II when $\phi = \frac{\pi}{2}$ then $\Delta \lambda = \frac{h}{m_0 c} = \lambda_c$

λ_c Compton wave length

$$\lambda_c = \frac{6.63 \times 10^{-34}}{9.1 \times 10^{-31} \times 3 \times 10^8}$$

$$\lambda_c = 0.0243 \text{ \AA}$$

The Compton wavelength is obtain in the direction of Normal to the incident photon.

Case-III when $\phi = \pi$ then $\Delta \lambda = \Delta \lambda_{\text{max}} = \frac{2h}{m_0 c}$

Ques. why Compton effect is not observed in visible light region? $\Delta \lambda_{\text{max}} = 0.0486 \text{ \AA}$

Ans. thus Compton effect is not observed in visible light.

Ques. why? because the change in wave length is negligible in comparison to the wave length of visible light

Que. Calculate the De-Broglie wavelength if the particle is n-electron and potential difference $V = 50$ volt e?

Ans.

$$d = \frac{h}{\sqrt{2mqV}} = \frac{12.28}{\sqrt{V}} \text{ \AA}$$

$$d = \frac{12.28}{\sqrt{50}} \text{ \AA}$$

$$d = 1.7366 \text{ \AA}$$

Que. Calculate D. broglie wavelength of a 15KV electron.

Soln -

$$d = \frac{h}{\sqrt{2mqV}}$$

$$E = 15 \text{ KeV} = 15 \times 10^3 \times 1.6 \times 10^{19} \text{ J}$$

$$d = \frac{6.63 \times 10^{-34}}{\sqrt{2 \times 9.1 \times 10^{-31} \times 15 \times 10^3 \times 1.6 \times 10^{19}}}$$

$$d = 0.1 \text{ \AA}$$

Que. Calculate the velocity and K.E. of a 1 Neutron having D.B.W 1 \text{ \AA}

Soln -

$$V = ?$$

$$E = ?$$

$$d = \frac{h}{mV}$$

$$V = \frac{h}{md} = \frac{6.63 \times 10^{-34}}{1.67 \times 10^{-27} \times 1 \times 10^{-10}}$$

$$V = 3.96 \times 10^3 \text{ m/sec}$$

$$d_n = \frac{0.286}{\sqrt{E_{eV}}} \text{ \AA}$$

$$\sqrt{E_{eV}} = \frac{0.286}{d} \times 10^{-10}$$

$$\sqrt{E_{eV}} = \frac{0.286}{1 \times 10^{-10}} \times 10^{-10}$$

$$E_{eV} = 0.082 \text{ eV}$$

Que. Calculate D.B.W. of an α particle exalted through a potential difference of 200 V.

Soln -

$$m_{\alpha} = 4m_n = 4 \times 1.67 \times 10^{-27} \text{ kg}$$

$$q_{\alpha} = 2e = 2 \times 1.6 \times 10^{-19} \text{ C}$$

$$d = \frac{h}{\sqrt{2mqV}} = \frac{6.63 \times 10^{-34}}{\sqrt{2 \times 4 \times 1.67 \times 10^{-27} \times 2 \times 1.6 \times 10^{-19} \times 200}}$$

$$d = 0.00716 \text{ \AA} \text{ OR } 716 \times 10^{12} \text{ m } 10^{19} \times 200$$

Que. Calculate the D.B.W. of a neutron having K.E. of 1 eV

Soln -

$$d_n = \frac{0.286}{\sqrt{E_{eV}}} \text{ \AA}$$

$$d_n = \frac{0.286}{\sqrt{1}}$$

$$d_n = 0.286 \text{ \AA}$$

Que. Can a photon and an electron of the same momentum have the same wavelength compared their wavelength in the λ have a same energy.

Solⁿ - 1 Given $p_{ph} = p_e$

$$dp_{ph} = \frac{h}{\lambda} \quad \text{--- (1)}$$

$$d_e = \frac{h}{p_e} \quad \text{--- (2)}$$

eqn / (1) & (2)

$$\frac{dp_{ph}}{d_e} = \frac{h}{dp_{ph}} \times \frac{p_e}{h} \quad \left\{ \begin{array}{l} p_{ph} = p_e \end{array} \right.$$

$$\frac{dp_{ph}}{d_e} = \frac{p_e}{dp_{ph}}$$

$$\boxed{dp_{ph} = d_e}$$

(ii) Given $E_{ph} = E_e = E$ (say)

$$d_e = \frac{h}{\sqrt{2mE}} \quad \text{--- (1)}$$

$$dp_{ph} = \frac{hc}{E_{ph}} \quad \text{--- (2)}$$

eg (1) / eq (2)

$$\frac{d_e}{dp_{ph}} = \frac{h}{\sqrt{2mE}} \times \frac{E}{hc} \quad \left\{ \begin{array}{l} E_{ph} = E_e = E \end{array} \right.$$

Que. Calculate D.B.W associated with Nitrogen. 30 atmospheric pressure and temp. 27°C and mass 4.65×10^{-26} kg.

Solⁿ - $T = 27^\circ C = 300K$

$$M_n = 4.65 \times 10^{-26} \text{ kg}$$

$$d = \frac{h}{\sqrt{3mKT}}$$

$$\left\{ k = 1.38 \times 10^{-23} \text{ J} \right.$$

$$d = \frac{6.63 \times 10^{-34}}{\sqrt{3 \times 4.65 \times 10^{-26} \times 1.38 \times 10^{-23} \times 300}}$$

$$\boxed{d = 0.275 \text{ \AA}} \quad \text{Ans.}$$

Imp

Que. N electron is bound in 1 dimensional potential well which has width 2.5×10^{-10} m assume the height of the well to be ∞ . Calculate the lowest two permitted energy values of e⁻.

$$E_n = \frac{n^2 h^2}{8mL^2}$$

$$E_1 = \frac{(6.63 \times 10^{-34})^2}{8 \times 9.1 \times 10^{-31} \times (2.5 \times 10^{-10})^2} \times n^2$$

$$E_n = 9.66 \times 10^{-19} \text{ n}^2 \text{ Joule}$$

Now

$$E_1 = 9.66 \times 10^{-19} \text{ Joule}$$

$$E_2 = 19.32 \times 10^{-19} \text{ Joule}$$

Que. A proton recoils back after striking an electron at rest. what is the change in wavelength of the photon. $\phi = \pi$

Ans.

$$\Delta\lambda = \frac{h}{m_0c} (1 - \cos\phi)$$

$$\Delta\lambda = \frac{6.63 \times 10^{-34}}{9.1 \times 10^{-31} \times 3 \times 10^8} (1 - (-1))$$

$$\Delta\lambda = 0.0486 \text{ \AA}$$

Ans.

Q. In a Compton experiment electrons of wavelength 0.015 \AA is scattered at 60° . find the wavelength scattered wavelength.

soln-

$$\phi = 60^\circ$$

$$\lambda' = ?$$

$$\lambda' - \lambda = \frac{h}{m_0c} (1 - \cos\phi)$$

$$\lambda' - 0.015 = \frac{6.63 \times 10^{-34}}{9.1 \times 10^{-31} \times 3 \times 10^8} (1 - \cos 60)$$

$$\lambda' - 0.015 = 0.0243 (1 - \frac{1}{2})$$

$$\lambda' = \frac{0.0243}{2} + 0.015$$

$$\lambda' = 0.027 \text{ \AA}$$

Que. An x-ray photon is found to have its wavelength doubled on being scattered to 90° . find the wavelength and energy of incident photon.

Ans.

$$\lambda = ?$$

$$E = ?$$

$$\lambda' - \lambda = \frac{h}{m_0c} (1 - \cos 90)$$

$$2\lambda - \lambda = 0.0243 (1 - 0) \quad \therefore \lambda' = 2\lambda$$

$$\lambda = 0.0243 \text{ \AA}$$

$$E = h\nu = hc/\lambda = \frac{6.63 \times 10^{-34} \times 3 \times 10^8}{0.0243 \times 10^{-10}}$$

$$E = 8.106 \times 10^{19} \text{ Joule}$$

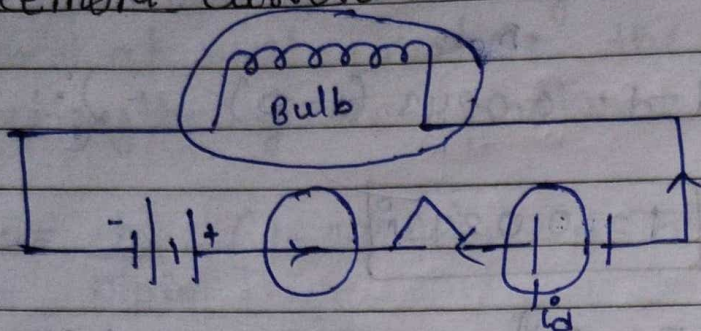
Questions

1. Define De broglie Hypothesis?
2. Define phase leave and group velocity?
3. Physical significance of wave function?
4. Define eigen value and eigen function?
5. what is modify and demodify radiation?
6. why Compton effect is not observed in visible region?

2 mark-7}

7. Derive different expression of De broglie wavelength?
8. (i) $v_g = v$ (ii) $v_p v_g = c^2$
9. Schrodinger time dependent and independent?
10. what is Compton effect?

② mark Displacement Current-



- Maxwell introduced that the magnetic field is produced by change in electric field.
- The current flowing due to the change in electric field b/w the plates of capacitor is called displacement current.

$$i_d = \epsilon_0 A \frac{\partial E}{\partial t}$$

$$i_d = A \frac{\partial D}{\partial t}$$

{ where $D = \epsilon_0 E$
 Electric displacement vector }

The magnetic field produced by displacement current is similar to the magnetic field produced by conventional current.

It produces as long as electric field is changing.

Gauss divergence theorem-

According to Gauss divergence theorem the surface integral may be changed to volume integral by taking the divergence of given vectors.

$$\text{i.e. } \oint_S \vec{A} \cdot d\vec{s} = \int_V \nabla \cdot \vec{A} dV$$

Stoke's theorem

According to Stoke's theorem the line integral may be change to surface integral by taking the curl of the given vector.

i.e.

$$\oint_C \vec{A} \cdot d\vec{l} = \int_S \nabla \times \vec{A} \cdot d\vec{s}$$

Equation of Continuity - The current is define as the Ratio of change to the time take.

Imp

$$\text{i.e. } i = \frac{q}{t}$$

current is also define as-

$$i = \oint_S \vec{J} \cdot d\vec{s} \quad \text{--- (1)}$$

$\therefore J =$ current density

For a small volume element the current is also define as weight rate of ^{negative} decrease of ~~two~~ charge with in that element.

$$\text{i.e. } i = - \frac{dQ}{dt} \quad \text{(2)}$$

The volume charge density is given by

$$\rho = \frac{Q}{V}$$

$$\text{the } Q = \int_V \rho dV$$

$$\text{the } i = - \frac{\partial}{\partial t} \int_V \rho dV \quad \text{(3)}$$

From eq (1) and (3)

$$\oint_S \vec{J} \cdot d\vec{s}' = \int_V \frac{\partial \rho}{\partial t} dV$$

Now the surface integral may be change to
Volume integral using Gauss divergence
Theorem -

then

$$\int_V \nabla \cdot \vec{J} \cdot dV = -\frac{\partial}{\partial t} \int_V \rho dV$$

then

$$\int_V \left(\nabla \cdot \vec{J} + \frac{\partial \rho}{\partial t} \right) dV = 0$$

then

$$\nabla \cdot \vec{J} = -\frac{\partial \rho}{\partial t} \quad (4)$$

$$\text{div } \vec{J} = -\frac{\partial \rho}{\partial t}$$

Eq (4) is known as eq of continuity.

Imp The physical significance of eqⁿ of continuity is that "The charge diverging from a small volume element is the current passing through the element."

physical significance of continuity eqⁿ-

Imp

The charge diverging from a small element is equal to the rate of decrease charge with that element.

Maxwell's Electromagnetic wave eqⁿ - All the fundamental law of electricity and magnetism are formulate in terms of Four equation which are known as Maxwell Em wave eqⁿ.

Maxwell also modified Amp. circuital law for static as well as dynamic charges.

(A) Maxwell Eqⁿ in Free Space -

$$\rho = 0 \quad \vec{J} = 0$$

$$(i) \quad \left. \begin{array}{l} \nabla \cdot \vec{E} = 0 \\ \nabla \cdot \vec{D} = 0 \end{array} \right\} \rightarrow (1) \quad \left\{ \because \vec{D} = \epsilon_0 \vec{E} \right\}$$

[Gauss Law of Electrostatic]

$$(ii) \quad \nabla \cdot \vec{B} = 0 \quad - (2) \quad \left\{ \because \vec{B} = \mu_0 \vec{H} \right\}$$

[Gauss Law of magnetostatic]

$$(iii) \quad \nabla \times \vec{E} = - \frac{\partial \vec{B}}{\partial t} \quad - (3) \quad \text{[Farade Law of Electro magnetic Induction]}$$

$$(iv) \quad \nabla \times \vec{H} = \epsilon_0 \frac{\partial \vec{E}}{\partial t} \quad - (4)$$

[Modify form of Amp. Circuital law]

(B) Maxwell Eqⁿ in differential form -

$$(i) \quad \left. \begin{array}{l} \nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0} \\ \nabla \cdot \vec{D} = \rho \end{array} \right\} \rightarrow (1)$$

$$(ii) \quad \nabla \cdot \vec{B} = 0 \quad - (2)$$

$$(iii) \quad \nabla \times \vec{E} = - \frac{\partial \vec{B}}{\partial t} \quad - (3)$$

Maxwell's Electromagnetic wave eqⁿ - All the fundamental law of electricity and magnetism are formulate in terms of four equation which are known as Maxwell Em wave eqⁿ.

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(A) Maxwell Eqⁿ in Free Space -

$$\rho = 0 \quad \vec{J} = 0$$

$$(i) \quad \left. \begin{array}{l} \nabla \cdot \vec{E} = 0 \\ \text{OR} \\ \nabla \cdot \vec{D} = 0 \end{array} \right\} \rightarrow (1) \quad \left\{ \begin{array}{l} \therefore \vec{D} = \epsilon_0 \vec{E} \end{array} \right\} \\ \text{[Gauss Law of Electrostatic]}$$

$$(ii) \quad \nabla \cdot \vec{B} = 0 \quad - (2) \quad \left\{ \begin{array}{l} \therefore \vec{B} = \mu_0 \vec{H} \end{array} \right\}$$

[Gauss Law of magnetostatic]

$$(iii) \quad \nabla \times \vec{E} = - \frac{\partial \vec{B}}{\partial t} \quad - (3) \quad \text{[Farade Law of Electro magnetic Induction]}$$

$$(iv) \quad \nabla \times \vec{H} = \epsilon_0 \frac{\partial \vec{E}}{\partial t} \quad - (4) \quad \text{[Modify form of Amp. Circuital Law]}$$

(B) Maxwell Eqⁿ in differential form -

$$(i) \quad \left. \begin{array}{l} \nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0} \\ \text{OR} \\ \nabla \cdot \vec{D} = \rho \end{array} \right\} \rightarrow (1)$$

$$(ii) \quad \nabla \cdot \vec{B} = 0 \quad - (2)$$

$$(iii) \quad \nabla \times \vec{E} = - \frac{\partial \vec{B}}{\partial t} \quad - (3)$$

$$(iv) \nabla \times \vec{H} = \vec{J} + \epsilon_0 \frac{\partial \vec{E}}{\partial t} \quad - (4) \quad \text{Imp}$$

(c) maxwell Eqⁿ in Integral form-

$$(i) \oint_S \vec{E} \cdot d\vec{s} = \frac{1}{\epsilon_0} \int_V \rho \, dv \quad - (1)$$

$$(ii) \oint_S \vec{B} \cdot d\vec{s} = 0 \quad - (2)$$

$$(iii) \oint_C \vec{E} \cdot d\vec{l} = - \int_S \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s} \quad - (3)$$

$$(iv) \oint_C \vec{H} \cdot d\vec{l} = \int_S \left(\vec{J} + \epsilon_0 \frac{\partial \vec{E}}{\partial t} \right) \cdot d\vec{s} \quad - (4)$$

Derivation of maxwell's in differential form-

(i) First eqⁿ \rightarrow

$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

According to Gauss Law of Electrostatic

$$\oint_S \vec{E} \cdot d\vec{s} = \frac{1}{\epsilon_0} \int_V \rho \, dv$$

using Gauss divergence theorem-

$$\int_V \nabla \cdot \vec{E} \cdot dV = \frac{1}{\epsilon_0} \int_V \rho \, dV$$

$$\int_V \left(\nabla \cdot \vec{E} - \frac{1}{\epsilon_0} \rho \right) dV = 0$$

$$\nabla \cdot \vec{E} - \frac{\rho}{\epsilon_0} = 0$$

$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

proved

(ii) Second eqⁿ → $\boxed{\nabla \cdot \vec{B}' = 0}$

According to Gauss law of Electromagnetic -
magnostatic -

$$\oint_S \vec{B}' \cdot d\vec{s}' = 0$$

Using Gauss divergence theorem -

$$\int_V \nabla \cdot \vec{B}' \, dV = 0$$

$$\boxed{\nabla \cdot \vec{B}' = 0}$$

proved

(iii) Third eqⁿ -

$$\boxed{\nabla \times \vec{E}' = - \frac{\partial \vec{B}'}{\partial t}}$$

The magnetic flux in term of magnetic flux density is define as

$$\phi = \oint_S \vec{B}' \cdot d\vec{s}' \quad \text{--- (1)}$$

According to Farade Law of Electromagnetic Induction -

$$e = - \frac{\partial \phi}{\partial t}$$

$$\boxed{e = - \oint_S \frac{\partial \vec{B}'}{\partial t} \cdot d\vec{s}'} \quad \text{--- (2)}$$

Let \vec{E} is the electric field of line element $d\vec{l}$ for a closed loop the EMF is define as the workdone to carry a unit charge or the closed loop that is-

$$\boxed{e = \oint_C \vec{E} \cdot d\vec{l}} \quad (3)$$

From eqⁿ 2 & (3)

$$\oint_C \vec{E} \cdot d\vec{l} = - \oint_S \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s}$$

using stoke's theorem-

$$\oint_C \nabla \times \vec{E} \cdot d\vec{s} = - \oint_S \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s}$$

$$\boxed{\nabla \times \vec{E} = - \frac{\partial \vec{B}}{\partial t}}$$

most
Imp

proved

(iv) Fourth equation -

or
Displacement
current

$$\nabla \times \vec{H} = \vec{J} + \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

Ampere's circuital law is define as-

$$\oint_C \vec{H} \cdot d\vec{l} = i \quad (1)$$

$$\oint_C \vec{H} \cdot d\vec{l} = \oint_S \vec{J} \cdot d\vec{s} \quad (2)$$

by using stoke's theorem-

$$\oint_C \nabla \times \vec{H}' \cdot d\vec{s}' = \oint_C \vec{J}' \cdot d\vec{s}'$$

$$\boxed{\nabla \times \vec{H}' = \vec{J}'} \quad - (3)$$

Taking divergence both sides.

$$\nabla \cdot \nabla \times \vec{H}' = \nabla \cdot \vec{J}'$$

$$\boxed{\nabla \cdot \vec{J}' = 0} \quad - (4) \quad \left\{ \nabla \cdot \nabla \times \vec{H}' = 0 \right\}$$

but, the eqⁿ according to eqⁿ of continuity-

$$\boxed{\nabla \cdot \vec{J}' = -\frac{\partial \rho}{\partial t}} \quad - (5)$$

If eqⁿ 4th and 5th is valid then $\rho = \text{const.}$

i.e. Amp. Circital Law is valid only

for static charge. Maxwell modified it for static as well as dynamic charges by adding a term -

\vec{J}'_d on R.H.S of eqⁿ (3).

$$\boxed{\nabla \times \vec{H}' = \vec{J}' + \vec{J}'_d} \quad - (6)$$

Taking divergence both sides

$$\nabla \cdot \nabla \times \vec{H}' = \nabla \cdot \vec{J}' + \nabla \cdot \vec{J}'_d$$

$$\nabla \cdot \vec{J}' + \nabla \cdot \vec{J}'_d = 0$$

$$\boxed{\nabla \cdot \vec{J}'_d = -\nabla \cdot \vec{J}'}$$

put value in eqⁿ (5)

$$\nabla \cdot \vec{J}_d = \frac{\partial \rho}{\partial t}$$

From Maxwell
1st eqⁿ $\nabla \cdot \vec{D} = \rho$

$$\nabla \cdot \vec{J}_d = \frac{\partial}{\partial t} (\nabla \cdot \vec{D})$$

$$\vec{J}_d = \frac{\partial \vec{D}}{\partial t}$$

$$\vec{J}_d = \epsilon_0 \frac{\partial \vec{E}}{\partial t} \quad \left\{ \vec{D} = \epsilon_0 \vec{E} \right\}$$

put in eqⁿ (6)

$$\boxed{\nabla \times \vec{H} = \vec{J} + \epsilon_0 \frac{\partial \vec{E}}{\partial t}}$$

Hence proved

Derivation of Maxwell's eqⁿ in integral form-

(i) First eqⁿ -

$$\boxed{\oint_S \vec{E} \cdot d\vec{s} = \frac{1}{\epsilon_0} \int_V \rho dV}$$

Maxwell first eqⁿ Differential form can be expressed as -

$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

Taking volume integral both side -

$$\int_V \nabla \cdot \vec{E} dV = \frac{1}{\epsilon_0} \int_V \rho dV$$

using Gauss divergence theorem -

$$\boxed{\oint_S \vec{E}' \cdot d\vec{s}' = \frac{1}{\epsilon_0} \int_V \rho dV}$$

proved

(ii) Second eqⁿ-

$$\oint_S \vec{B}' \cdot d\vec{s}' = 0$$

maxwell second eqⁿ in differential form -

$$\nabla \cdot \vec{B}' = 0$$

Taking volume integral

$$\int_V \nabla \cdot \vec{B}' \cdot dV = 0$$

Using Gauss divergence theorem-

$$\boxed{\oint_S \vec{B}' \cdot d\vec{s}' = 0}$$

(iii) Third eqⁿ-

$$\boxed{\oint_C \vec{E}' \cdot d\vec{l}' = - \int \frac{\partial \vec{B}'}{\partial t} \cdot d\vec{s}'}$$

maxwell third eqⁿ in differential form is defined as-

$$\nabla \times \vec{E}' = - \frac{\partial \vec{B}'}{\partial t}$$

Taking surface integral both sides

$$\oint_S \nabla \cdot \vec{E}' \cdot d\vec{s}' = \oint_S - \frac{\partial \vec{B}'}{\partial t} \cdot d\vec{s}'$$

by using Stokes's theorem -

$$\oint_C \vec{E}' \cdot d\vec{l}' = - \oint_S \frac{\partial \vec{B}'}{\partial t} \cdot d\vec{s}$$

proved

Fourth eqⁿ -

$$\oint_C \vec{H}' \cdot d\vec{l}' = \oint_S \left(\vec{J}' + \epsilon_0 \frac{\partial \vec{E}'}{\partial t} \right) \cdot d\vec{s}$$

Maxwell Fourth eqⁿ in differential form is defined as -

$$\nabla \times \vec{H}' = \vec{J}' + \epsilon_0 \frac{\partial \vec{E}'}{\partial t}$$

Taking surface integral -

$$\oint_S \nabla \times \vec{H}' \cdot d\vec{s}' = \oint_S \left(\vec{J}' + \epsilon_0 \frac{\partial \vec{E}'}{\partial t} \right) \cdot d\vec{s}'$$

by Stokes's theorem -

$$\oint_C \vec{H}' \cdot d\vec{l}' = \oint_S \left(\vec{J}' + \epsilon_0 \frac{\partial \vec{E}'}{\partial t} \right) \cdot d\vec{s}'$$

proved

Physical significance of Maxwell's eqⁿ -

First eqⁿ - To introduce that total Electric flux and closed by Hypothetical gaussian surface is $\frac{1}{\epsilon_0}$ times the total charge

and closed within that surface.

Thus, it represents Gauss Law of electrostatics.

② Second eqⁿ-

It introduces that magnetic monopoles don't exist. So it represents Gauss Law of magnetostatics.

③ Third eqⁿ-

It introduces that the rate of change of magnetic flux produces an induced emf. Hence, it represents Faraday Law of Electro-magnetic Induction.

④ Fourth eqⁿ- It states that the magnetic field is produced by conventional current as well as change in electric field. It represents modified form of Ampere's circuital law.

Very Imp

Poynting vector-

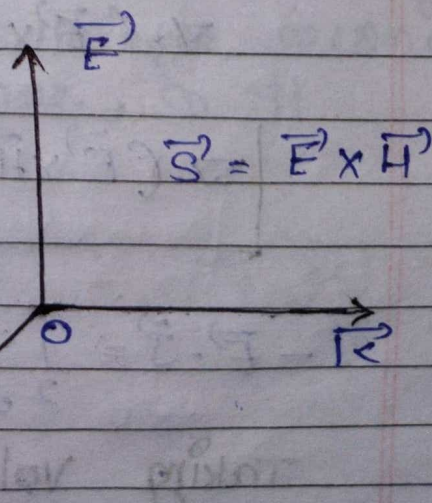
Numerical - 7

Derivation - 7

Short - 2

The amount of energy passing per unit area per second perpendicular to

$E-H$ plane is known as Poynting vector.



i.e. $\vec{S} = \vec{E} \times \vec{H}$
 $\vec{S} = EH \sin \theta \hat{n}$

$$|\vec{S}| = EH \quad \text{Joule/sec.} \cdot \text{m}^2$$

Poynting theorem -
 Maxwell third and 4th eq in differential
 form can be expressed as -

$$\nabla \times \vec{E} = -\mu \frac{\partial \vec{H}}{\partial t} \quad (1)$$

$$\nabla \times \vec{H} = \epsilon \frac{\partial \vec{E}}{\partial t} + \vec{J} \quad (2)$$

Taking dot product by \vec{H} and \vec{E} with
 eq (1) & (2) respectively -

$$\vec{H} \cdot \nabla \times \vec{E} = -\mu \vec{H} \cdot \frac{\partial \vec{H}}{\partial t} \quad (3)$$

$$\vec{E} \cdot \nabla \times \vec{H} = \epsilon \vec{E} \cdot \frac{\partial \vec{E}}{\partial t} + \vec{E} \cdot \vec{J} \quad (4)$$

we know that -

$$\nabla \cdot (\vec{E} \times \vec{H}) = \vec{H} \cdot \nabla \times \vec{E} - \vec{E} \cdot \nabla \times \vec{H} \quad (5)$$

$$\nabla \cdot (\vec{E} \times \vec{H}) = -\mu \vec{H} \cdot \frac{\partial \vec{H}}{\partial t} - \epsilon \vec{E} \cdot \frac{\partial \vec{E}}{\partial t} - \vec{E} \cdot \vec{J}$$

$$\nabla \cdot (\vec{E} \times \vec{H}) = -\frac{1}{2} \frac{\partial}{\partial t} (\mu \vec{H} \cdot \vec{H} + \epsilon \vec{E} \cdot \vec{E}) - \vec{E} \cdot \vec{J}$$

(6)

$$-\vec{E} \cdot \vec{J} = \frac{1}{2} \frac{\partial}{\partial t} (\vec{B} \cdot \vec{H} + \vec{D} \cdot \vec{E}) + \nabla \cdot (\vec{E} \times \vec{H})$$

Taking volume Integral on both sides -

$$-\int_V \vec{E} \cdot \vec{J} \, dV = \frac{1}{2} \frac{\partial}{\partial t} \int_V (\vec{B} \cdot \vec{H} + \vec{D} \cdot \vec{E}) \, dV +$$

$$\int_V \nabla \cdot (\vec{E} \times \vec{H}) \, dV \rightarrow \text{Gauss div. theorem}$$

$$-\int_V \vec{E} \cdot \vec{J} \, dV = \underbrace{\frac{1}{2} \frac{\partial}{\partial t} \int_V (\vec{B} \cdot \vec{H} + \vec{D} \cdot \vec{E}) \, dV}_{\text{1st term}} + \underbrace{\oint_S (\vec{E} \times \vec{H}) \cdot d\vec{S}}_{\text{third term}} \quad \text{second term} \quad \text{⑦}$$

This is known as Poynting theorem or Conservation law of energy in electro magnetic wave.

The physical significance of Poynting theorem-

- (I) 1st term- It represents the transmission of energy due to the motion of charge.
- (II) 2nd term- It represents the sum of magnetic and electric potential energy's.
- (III) 3rd term- It represents the amount of energy crossing per unit area per second perpendicular to $\vec{E} \times \vec{H}$ plane, so it represents Poynting vector.

$$\text{Poynting vector} = \vec{E} \times \vec{H} = \vec{S}$$

V. Imp # Propagation of EM wave in free space-

maxwell Electromagnetic wave eqⁿ in free space can be described as-

$$\nabla \cdot \vec{E} = 0 \quad \text{--- (1)}$$

$$\nabla \cdot \vec{H} = 0 \quad \text{--- (2)}$$

$$\nabla \times \vec{E} = -\mu_0 \frac{\partial \vec{H}}{\partial t} \quad \text{--- (3)}$$

$$\nabla \times \vec{H} = \epsilon_0 \frac{\partial \vec{E}}{\partial t} \quad \text{--- (4)}$$

Taking curl both sides in eqⁿ (3)

$$\nabla \times \nabla \times \vec{E} = -\mu_0 \frac{\partial}{\partial t} \nabla \times \vec{H}$$

$$\boxed{\nabla \times \nabla \times \vec{E} = -\mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}} \quad \left\{ \text{using eqⁿ (4)} \right\} \quad \text{--- (5)}$$

We know that-

$$\boxed{\vec{A} \times \vec{B} \times \vec{C} = \vec{B}(\vec{A} \cdot \vec{C}) - \vec{C}(\vec{A} \cdot \vec{B})} \quad \text{--- (6)}$$

by comparing eqⁿ (5) and (6)

$$\boxed{\nabla \times \nabla \times \vec{E} = \nabla(\nabla \cdot \vec{E}) - (\nabla \cdot \nabla)\vec{E}} \quad \left\{ \text{eqⁿ (1) using } \right.$$

$$\boxed{\nabla \times \nabla \times \vec{E} = 0 - \nabla^2 \vec{E}} \quad \left. \text{put eqⁿ (5)} \right.$$

$$-\nabla^2 \vec{E} = -\mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\boxed{\nabla^2 \vec{E} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}} \quad \text{--- (7)}$$

similarly on taking the curl on (eq) (4) both side we have -

$$\nabla^2 \vec{H} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{H}}{\partial t^2} \quad (8)$$

we obtain two similar electromagnetic wave equation \vec{E} and another \vec{H} for this the general equation of EM wave is given by -

$$\nabla^2 \phi = \frac{1}{v^2} \frac{\partial^2 \phi}{\partial t^2} \quad (9)$$

Compare eq (7), (8), (9)

$$\frac{1}{v^2} = \mu_0 \epsilon_0$$

$$v = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$$

Now,

$$\mu_0 = 4\pi \times 10^{-7}$$

$$\epsilon_0 = \frac{1}{4\pi \times 9 \times 10^9}$$

$$v = \frac{1}{\sqrt{4\pi \times 10^{-7} \times 4\pi \times 9 \times 10^9}} = \frac{1}{\sqrt{4\pi \times 10^{-7} \times 1}} = \frac{1}{\sqrt{4\pi \times 9 \times 10^9}} = 3 \times 10^8 = c \text{ m/s}$$

$$v = \frac{1}{(10^8/3)} = 3 \times 10^8 = c \text{ m/s}$$

$$v = c = 3 \times 10^8 \text{ m/s}$$

Thus, electromagnetic wave travel with the speed of light in free space (vacuum).

Transverse nature of EM wave -

The Electromagnetic wave

Eq^y for \vec{E} and \vec{H} are

given by -

$$\nabla^2 \vec{E} - \frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = 0 \quad \text{--- (1)}$$

$$\nabla^2 \vec{H} - \frac{1}{c^2} \frac{\partial^2 \vec{H}}{\partial t^2} = 0 \quad \text{--- (2)}$$

The sol^y of eq^y (1) & eq^y (2) are given by

$$\vec{E} = \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} \quad \text{--- (3)}$$

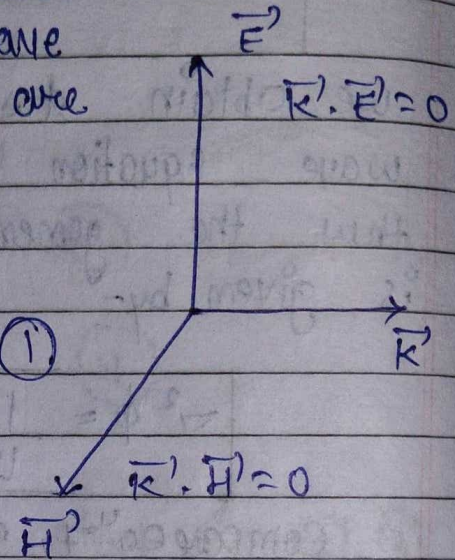
$$\vec{H} = \vec{H}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} \quad \text{--- (4)}$$

The maxwell Ist and IInd eq^y in Free space are defined as -

$$\nabla \cdot \vec{E} = 0 \quad \text{--- (5)}$$

$$\nabla \cdot \vec{H} = 0 \quad \text{--- (6)}$$

put value in eq^y (5) of \vec{E}



$$\nabla \cdot \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} = 0$$

$$i \vec{k} \cdot \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} = 0$$

↓
 \vec{E}

$$i \vec{k} \cdot \vec{E} = 0$$

$$\boxed{\vec{k} \cdot \vec{E} = 0} \quad - (7)$$

$$\vec{k} \perp \vec{E}$$

From eq^y 4th and 6th

$$\nabla \cdot \vec{H}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} = 0$$

$$i \vec{k} \cdot \vec{H} = 0$$

$$\boxed{\vec{k} \cdot \vec{H} = 0} \quad - (8)$$

$$\vec{k} \perp \vec{H}$$

From eq^y (7) and eq^y (8) it is clear that k is perpendicular to both \vec{E} and \vec{H} thus electromagnetic wave is transverse in nature.

$$\nabla = \frac{\partial}{\partial x} \hat{i} + \frac{\partial}{\partial y} \hat{j} + \frac{\partial}{\partial z} \hat{k}$$

$$\vec{E}_0 = E_0 x \hat{i} + E_0 y \hat{j} + E_0 z \hat{k}$$

$$\vec{k} = k_x \hat{i} + k_y \hat{j} + k_z \hat{k}$$

$$\vec{r} = x \hat{i} + y \hat{j} + z \hat{k}$$

$$\vec{k} \cdot \vec{r} = x k_x + y k_y + z k_z$$

$$\vec{k} \cdot \vec{E}_0 = k_x E_0 x + k_y E_0 y + k_z E_0 z$$

Now

$$\left(\frac{\partial}{\partial x} \hat{i} + \frac{\partial}{\partial y} \hat{j} + \frac{\partial}{\partial z} \hat{k} \right) \cdot (E_0 x \hat{i} + E_0 y \hat{j} + E_0 z \hat{k}) e^{i(k_x x + k_y y + k_z z - \omega t)} = 0$$

$$E_0 x \frac{\partial}{\partial x} e^{i(k_x x + k_y y + k_z z - \omega t)} + E_0 y \frac{\partial}{\partial y} e^{i(k_x x + k_y y + k_z z - \omega t)} + E_0 z \frac{\partial}{\partial z} e^{i(k_x x + k_y y + k_z z - \omega t)} = 0$$

$$(i E_0 k_x x + i E_0 y k_y + i E_0 z k_z) e^{i(\vec{k} \cdot \vec{r} - \omega t)} = 0$$

$$i (k_x E_0 x + k_y E_0 y + k_z E_0 z) e^{i(\vec{k} \cdot \vec{r} - \omega t)} = 0$$

$$i \vec{k} \cdot \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} = 0$$

$$i \vec{k} \cdot \vec{E}' = 0$$

$$\boxed{\vec{k} \cdot \vec{E}' = 0}$$

- (7)

$$\boxed{\vec{k} \perp \vec{E}'}$$

Similarly

$$\nabla \cdot \vec{H}' e^{i(\vec{k} \cdot \vec{r} - \omega t)} = 0$$

$$i \vec{k} \cdot \vec{H}' = 0$$

$$\boxed{\vec{k} \cdot \vec{H}' = 0}$$

- (8)

$$\boxed{\vec{k} \perp \vec{H}'}$$

The Maxwell IIIrd and IVth eq^s can be expressed as -

$$\nabla \times \vec{E} = -\mu_0 \frac{\partial \vec{H}}{\partial t} \quad \text{--- (9)}$$

$$\nabla \times \vec{H} = \epsilon_0 \frac{\partial \vec{E}}{\partial t} \quad \text{--- (10)}$$

From eq^y (3), (9) & (10) we have -

$$\nabla \times \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} = -\mu_0 \frac{\partial}{\partial t} \vec{H}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

$$i \vec{k} \times \vec{E} = -i \omega \mu_0 \vec{H}$$

$$\boxed{\vec{k} \times \vec{E} = \omega \mu_0 \vec{H}} \quad \text{--- (11)}$$

From eq^y (3), (4) & (10) -

$$\nabla \times \vec{H}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} = \epsilon_0 \frac{\partial}{\partial t} \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

$$i \vec{k} \times \vec{H} = -i \omega \epsilon_0 \vec{E}$$

$$\boxed{\vec{k} \times \vec{H} = -\omega \epsilon_0 \vec{E}} \quad \text{--- (12)}$$

From eq^y (7), (8), (11) & (12) - It is clear that \vec{E} , \vec{H} and \vec{k} are mutually perpendicular.

From eq^y (11) we have -

$$k \times \vec{E} = \mu_0 \omega \vec{H}$$

$$k \hat{n} \times \vec{E} = \mu_0 \omega \vec{H}$$

$$k |\hat{n} \times \vec{E}| = \mu_0 \omega |\vec{H}|$$

$$k |\vec{E}| = \mu_0 \omega |\vec{H}|$$

$$\frac{|\vec{E}|}{|\vec{H}|} = \frac{\mu_0 \omega}{k} = \mu_0 c \left\{ \frac{\omega}{k} = c \right\}$$

$$\frac{|\vec{E}|}{|\vec{H}|} = \frac{\mu_0}{\sqrt{\mu_0 \epsilon_0}} = \sqrt{\frac{\mu}{\epsilon_0}}$$

$$\frac{|\vec{E}|}{|\vec{H}|} = \frac{4\pi \times 10^{-7}}{\sqrt{4\pi \times 9 \times 10^9}}$$

$$\frac{|\vec{E}|}{|\vec{H}|} = 376.72 \text{ ohm}$$

$$\boxed{\frac{E}{H} = 376.72 \ \Omega}$$

Imp propagation of EM wave in conducting medium.
 In conducting medium the charge volume density is zero and current density define by ohm's law.

i.e. $\rho = 0$, $\vec{J} = \sigma \vec{E}$

then maxwell Eq^y in conducting medium is define as-

$$\nabla \cdot \vec{E}' = 0 \quad (1)$$

$$\nabla \cdot \vec{H}' = 0 \quad (2)$$

$$\nabla \times \vec{E}' = -\mu \frac{\partial \vec{H}'}{\partial t} \quad (3)$$

$$\nabla \times \vec{H}' = \epsilon_0 \frac{\partial \vec{E}'}{\partial t} + \sigma \vec{E}' \quad (4)$$

Taking curl both side in eq^y (3)

$$\nabla \times \nabla \times \vec{E}' = -\mu \frac{\partial}{\partial t} \nabla \times \vec{H}'$$

$$\nabla \times \nabla \times \vec{E} = -\mu \frac{\partial}{\partial t} (\sigma \vec{E} + \epsilon \frac{\partial \vec{E}}{\partial t})$$

$$\left\{ \vec{A} \times (\vec{B} \times \vec{C}) = \vec{B}(\vec{A} \cdot \vec{C}) - \vec{C}(\vec{A} \cdot \vec{B}) \right\}$$

$$\nabla(\nabla \cdot \vec{E}) - \nabla^2 \vec{E} = -\mu \sigma \frac{\partial \vec{E}}{\partial t} + \mu \epsilon \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\boxed{\nabla^2 \vec{E} - \mu \sigma \frac{\partial \vec{E}}{\partial t} - \mu \epsilon \frac{\partial^2 \vec{E}}{\partial t^2} = 0} \quad \text{--- (5)}$$

similarly taking curl in eq (4) both sides -

$$\boxed{\nabla^2 \vec{H} - \mu \sigma \frac{\partial \vec{H}}{\partial t} - \mu \epsilon \frac{\partial^2 \vec{H}}{\partial t^2} = 0} \quad \text{--- (6)}$$

The solution of eq (5) & eq (6) are given by -

$$\vec{E} = \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} \quad \text{--- (7)}$$

$$\vec{H} = \vec{H}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} \quad \text{--- (8)}$$

From eq (5) & (7) we get -

$$\nabla^2 \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} - \mu \sigma \frac{\partial}{\partial t} \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} - \mu \epsilon \frac{\partial^2}{\partial t^2} \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} = 0$$

$$[-k^2 + i\omega\mu\sigma + \omega^2\mu\epsilon] \vec{E} = 0$$

$$-k^2 + i\omega\mu\sigma + \omega^2\mu\epsilon = 0$$

$$k^2 = \omega^2\mu\epsilon + i\omega\mu\sigma \quad \text{--- (9)}$$

$$\boxed{k^2 = \omega \left[\mu \epsilon \left(1 + \frac{i\sigma}{\omega \epsilon} \right) \right]} \quad \text{--- (10)}$$

Let k is a complex junction say -

$$k = A + iB \quad \text{--- (11)}$$

$$k^2 = A^2 - B^2 + 2iAB \quad \text{--- (12)}$$

From eqⁿ (11) & (12) we get -

$$A^2 - B^2 = \omega^2 \mu \epsilon \quad \text{--- (13)}$$

$$2AB = \mu \omega \sigma \quad \text{--- (14)}$$

The sol^y eq^y (13) & (14) is given by -

$$A = \omega \sqrt{\mu \epsilon} \left[\frac{\sqrt{1 + \left(\frac{\sigma}{\omega \epsilon}\right)^2} + 1}{2} \right]^{1/2} \quad \text{--- (15)}$$

$$B = \omega \sqrt{\mu \epsilon} \left[\frac{\sqrt{1 + \left(\frac{\sigma}{\omega \epsilon}\right)^2} - 1}{2} \right]^{1/2} \quad \text{--- (16)}$$

For good conductor $\frac{\sigma}{\omega \epsilon} \gg \gg \gg 1$

$$A \approx \omega \sqrt{\mu \epsilon} \left[\frac{\frac{\sigma}{\omega \epsilon}}{2} \right]^{1/2} = B$$

$$A \approx \omega \sqrt{\mu \epsilon} \left[\frac{\sigma}{2\omega \epsilon} \right]^{1/2}$$

$$A \approx \left[\frac{\omega \sigma \mu}{2} \right]^{1/2} = B \quad \text{--- (17)}$$

then

$$\vec{E} = \vec{E}_0 e^{-B \cdot \vec{x}} e^{i(A \cdot \vec{x} - \omega t)} \quad \text{--- (18)}$$

$$\vec{H} = \vec{H}_0 e^{-B \cdot \vec{x}} e^{i(A \cdot \vec{x} - \omega t)} \quad \text{--- (19)}$$

The velocity of E.M. wave in conducting medium is given by-

$$v = \frac{\omega}{k} = \frac{\omega}{A}$$

$$v = \frac{\omega}{\omega \sqrt{\mu \epsilon}} \left(\frac{1 + \left(\frac{\sigma}{\omega \epsilon} \right)^2 + 1}{2} \right)^{-1/2}$$

$$v = \frac{1}{\sqrt{\mu \epsilon}} \left(\frac{1 + \left(\frac{\sigma}{\omega \epsilon} \right)^2 + 1}{2} \right)^{-1/2} \quad \text{--- (20)}$$

2-mark
#

Skin depth and skin effect - The distance (depth) travel by electromagnetic wave in conducting medium to reduce its amplitude $\frac{1}{e}$ time the amplitude

at the surface of the conductor is called skin depth.

$$\delta = \frac{1}{B} = \left(\frac{2}{\omega \mu} \right)^{1/2}$$

From here, it is clear that skin depth will be decreases on increasing frequency this is known as skin effect. Thus the long transmission of electricity take place at high frequency to minimize the power loss.

V. Imp
Q.

If the earth receives 2 cal/min cm^2 of solar energy what are the amplitude of electric and magnetic fields of radiations.

Soln-

$$\vec{S} = \vec{E} \times \vec{H}$$

$$|\vec{S}| = EH = \frac{2 \times 4.2}{60 \times 10^4}$$

$$EH = 1400 \quad \text{--- (1)}$$

We know that -

$$\frac{E}{H} = 376.72 \quad \text{--- (2)}$$

eq (1) \times eq (2)

$$EH \times \frac{E}{H} = 1400 \times 376.72 = 527408$$

$$E = 726.22 \text{ V/m}$$

Now,

$$E_0 = \sqrt{2} E$$

$$E_0 = \sqrt{2} \times 726.22$$

$$E_0 = 1027.03 \text{ V/m}$$

Now From eq (2)

$$\frac{E}{H} = 376.72$$

$$H = \frac{E}{376.72} = \frac{726.22}{376.72}$$

$$H = 1.92 \text{ AT/m}$$

Now,

$$H_0 = \sqrt{2} H$$

$$H_0 = \sqrt{2} \times 1.92 = 2.71 \text{ AT/m}$$

Assuming that all the energy from a 1000 watt lamp is mediated uniformly calculate the avg. values of the intensity of electric and magnetic field of radiation at the distance 2m from lamp.

$$|\vec{S}| = EH = \frac{1000}{4\pi} \text{ J/ sec-m}^2$$

$$EH = \frac{1000}{4\pi(2)^2} = \frac{1000}{16\pi}$$

$$EH = 19.89 \quad \text{--- (1)}$$

$$\frac{E}{H} = 376.72 \quad \text{--- (2)}$$

$$\text{Eq (1)} \times \text{Eq (2)}$$

$$EH \times \frac{E}{H} = 19.89 \times 376.72$$

$$E^2 = 7492.9608$$

$$E = 86.59 \text{ V/m}$$

From eq (2) -

$$\frac{E}{H} = 376.72$$

$$H = \frac{E}{376.72} = \frac{86.59}{376.72}$$

$$H = 0.23 \text{ AT/m}$$

Que. The sun light is strikes the upper atmosphere of earth with energy flux $1.38 \text{ kwatt m}^{-2}$, what will be the peak value of electric and magnetic fields at the point.

Ans.

$$\vec{S} = \vec{E} \times \vec{H}$$

$$|\vec{S}| = EH = 1.38 \times 10^3 \text{ J sec}^{-1} \text{ m}^{-2}$$

$$EH = 1.38 \times 10^3 \quad \text{--- (1)}$$

We know that

$$\frac{E}{H} = 376.72 \quad \text{--- (2)}$$

$$\text{Eq (1)} \times \text{Eq (2)}$$

$$E^2 = 1.38 \times 10^3 \times 376.72$$

$$E^2 = 519873.6$$

$$E = 721.022 \text{ N/m}$$

$$E_0 = \sqrt{2} E$$

$$E_0 = \sqrt{2} \times 721.022$$

$$E_0 = 1019.68 \text{ N/m}$$

From eq (2) we get-

$$H = \frac{E}{376.72} = \frac{721.022}{376.72}$$

$$H = 1.91 \text{ AT/m}$$

$$H_0 = \sqrt{2} H = \sqrt{2} \times 1.91 = 2.705 \text{ AT/m}$$

$$B_0 = \mu_0 H_0$$

$$B_0 = 4\pi \times 10^{-7} \times 2.705$$

$$B = 3.4 \times 10^{-6} \text{ wb/m}^2$$

Que. The max. Electric field in a plane electromagnetic wave 10^2 N/C . The wave is going in the x -direction and Electric field in y -direction find max. magnetic field in wave and its direction.

Ans.

$$\frac{E_0}{H_0} = \mu_0 c$$

$$\frac{E_0}{\mu_0 H_0} = c$$

$$\frac{E_0}{B_0} = c$$

$$B_0 = \frac{E_0}{c} = \frac{10^2}{3 \times 10^8} = 3.33 \times 10^{-7} \text{ Tesla}$$

Along z -direction

Que. Using Maxwell eqⁿ $\text{Curl } \vec{B} = \mu_0 \left(\vec{J} + \frac{\partial \vec{D}}{\partial t} \right)$
prove that $\text{div } \vec{D} = \rho$

Ans. Taking div of both sides -

$$\nabla \cdot \nabla \times \vec{B} = \mu_0 \left(\nabla \cdot \vec{J} + \frac{\partial}{\partial t} \nabla \cdot \vec{D} \right)$$

$$0 = \mu_0 \left(\nabla \cdot \vec{J} + \frac{\partial}{\partial t} \nabla \cdot \vec{D} \right)$$

$$\frac{\partial}{\partial t} \nabla \cdot \vec{D} = - \nabla \cdot \vec{J}$$

$$\left\{ \nabla \cdot \vec{J} = -\frac{\partial \rho}{\partial t} \right\}$$

$$\frac{\partial}{\partial t} (\nabla \cdot \vec{D}) = \frac{\partial \rho}{\partial t}$$

$$\boxed{\nabla \cdot \vec{D} = \rho}$$

proved

Calculate the skin depth of at the Frequency
71.6 MHz in Aluminium the related
parameter for aluminium are $\mu = \mu_0$
 $\mu = \mu_0 = 4\pi \times 10^{-7} \text{ N/Amp}^2$
 $\sigma = 3.54 \times 10^7 \text{ S/m}$

We know that -

$$\delta = \left(\frac{2}{\omega \sigma \mu} \right)^{1/2}$$

$$\left[\mu \sigma = \frac{\mu}{\mu_0} \right]$$

$$\delta = \left(\frac{2}{2\pi f \sigma \mu} \right)^{1/2}$$

$$\delta = \left(\frac{1}{3.14 \times 71.6 \times 10^6 \times 3.54 \times 10^7 \times 4 \times 3.14 \times 10^{-7}} \right)^{1/2}$$

$$\delta = 10 \times 10^{-6} \text{ m}$$

$$\boxed{\delta = 10 \mu\text{m}}$$

UNIT-3

Wave Optics

Page No.

Date

Coherent source - The light sources of same intensity, same wavelength, exactly same amplitude and zero(0) or constant phase difference are known as coherent sources.

Two independent light sources never be coherent \rightarrow

Light is produced due to the deexcitation of billion of atoms from excited state to lower energy state by emitting photon. The phase of the emitted photon change randomly at a point so the phase difference at any point is not zero(0) or const. Thus, two independent light source never be coherent.

Interference - When light from two or more than two coherent light sources propagate along the same direction then they superimpose constructively at certain point which are known as constructive interference and at other places they interfere destructively which are known as destructive interference.

The modification of in the intensity of resultant light when two coherent light propagate along the same direction is called interference.

Condition for sustain interference - The condition for sustain interference are as follows -

1. The intensity of light source must be same.
2. The wavelength of the source should be same.
- The amplitude of the light λ should be same.
- The phase difference b/w sources should be zero or 2π .
- The sources should be narrow.
- The light from the sources must be propagate along the same direction.
- The separation b/w the sources should be smaller.
- The distance b/w the sources and screen should be larger.
- If the interference occurs b/w the polarized light then the plane of polarization must be same.

Super position principle - According to super position principle the displacement of resultant light is the sum of displacement due to individual wave:

$$y = y_1 + y_2$$

$$y_1 = a_1 \sin \omega t$$

$$y_2 = a_2 \sin(\omega t + \phi) \quad \text{--- (1)}$$

$$y = a_1 \sin \omega t + a_2 \sin(\omega t + \phi) \quad \text{--- (2)}$$

$$\{ \sin(A+B) = \sin A \cos B + \cos A \sin B \}$$

$$y = a_1 \sin \omega t + a_2 \sin \omega t \cos \phi + a_2 \sin \phi \cos \omega t$$

$$y = (a_1 + a_2 \cos \phi) \sin \omega t + a_2 \sin \phi \cos \omega t \quad \text{--- (3)}$$

$$\text{Let } a_1 + a_2 \cos \phi = A \cos \delta \quad \text{--- (4)}$$

$$a_2 \sin \phi = A \sin \delta \quad \text{--- (5)}$$

$$\text{From eq 4, 5, 4, 5} \quad \text{--- (6)}$$

$$y = A \cos \delta \sin \omega t + A \sin \delta \cos \omega t$$

$$\text{Now, eq 4}^2 + \text{eq 5}^2 \quad \text{--- (7)}$$

$$A^2 = a_1^2 + a_2^2 + 2 a_1 a_2 \cos \phi \quad \text{--- (7)}$$

The intensity of wave is directly proportional to the square of amplitude.

$$I \propto A^2$$

$$I = A^2$$

$$I = a_1^2 + a_2^2 + 2 a_1 a_2 \cos \phi \quad \text{--- (8)}$$

Condition of maxima - The intensity will be maxima when -

$$\cos \phi = 1$$

$$\text{cos } \phi = \cos(2n\pi)$$

$$\text{so } \phi = 2n\pi$$

$$\text{path diff } (\Delta) = \frac{d}{2\lambda} \text{ phase diff.}$$

$$n = 0, 1, 2, \dots$$

$$\Delta = \frac{d}{2\lambda} \times 2n\pi$$

$$\Delta = n\lambda \quad \text{max.} \quad n=0, 1, 2, \dots$$

$$I_{\text{max}} \approx q_1^2 + q_2^2 + 2q_1q_2 \quad \text{--- (9)}$$

$$I_{\text{max}} \approx (q_1 + q_2)^2 \quad \text{--- (10)}$$

For minima - The intensity will be minimum when

$$\cos \phi = -1$$

$$\cos \phi = \cos (2n+1)\pi$$

$$\phi = (2n+1)\pi \quad n=0, 1, 2, \dots$$

then path diff.

$$\Delta = \frac{d}{2\lambda} \times (2n+1)\pi$$

$$\Delta = (2n+1) \frac{d}{2} \quad \text{minima}$$

$$I_{\text{min}} = q_1^2 + q_2^2 - 2q_1q_2 \quad \text{--- (11)}$$

$$I_{\text{min}} = (q_1 - q_2)^2 \quad \text{--- (12)}$$

Interference and conservation laws of energy - The avg. intensity of interference pattern can be expressed as

$$I_{\text{av}} = \frac{I_{\text{max}} + I_{\text{min}}}{2}$$

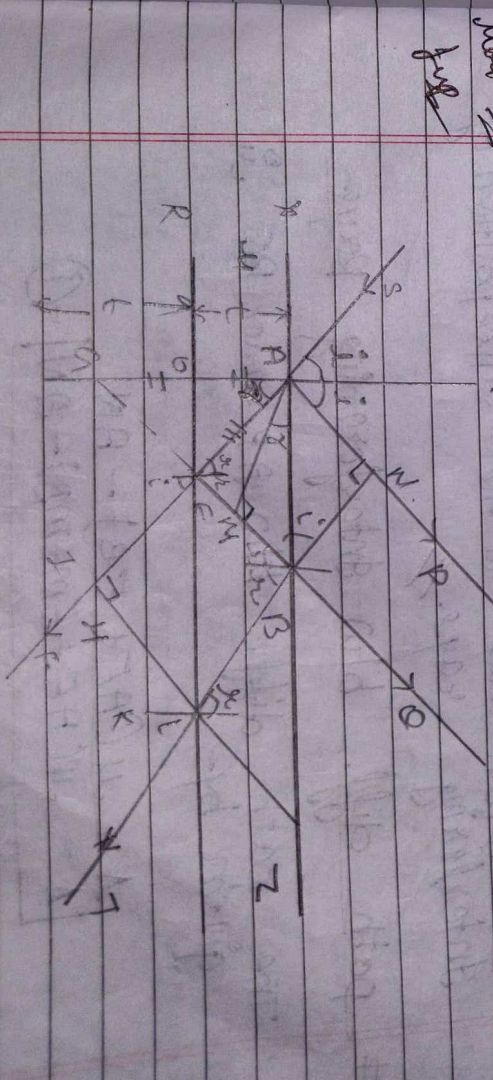
$$I_{\text{av}} = \frac{q_1^2 + q_2^2 + 2q_1q_2 + q_1^2 + q_2^2 - 2q_1q_2}{2}$$

$$I_{\text{av}} = \frac{2q_1^2 + 2q_2^2}{2}$$

$$I_{\text{av}} = q_1^2 + q_2^2$$

Thus, the avg. intensity in the sum of the intensity of individual wave. So, interference pattern is consistent with conservation law of energy. Whatever energy may be apparently present at maxima, at minima actually present at maxima.

Interference due to uniform thin films - Parallel the film -



When monochromatic light incident on the upper surface of uniform thin film then interference of path due to the superposition of interfering rays.

- (I) In Reflected system
(II) In Transmitted system

Due to the division of amplitude,

lets consider a thin film of refractive index (μ) and thickness (t) is bounded between the upper surface ($X-Y$) and lower surface ($R-Z$).

When monochromatic light incident at a on the upper surface then it suffer multiple reflection and refraction from upper and lower surface of the film and interfered pattern is obtain due to the superposition of interfering rays.

path diff. b/w interfering rays-

The path diff. b/w AP and BO is given by-

$$A = \mu(\mu t \cos r - \mu n)$$

$$\Delta = \mu(\mu t \cos r - \mu n) - \mu n \quad \text{--- (1)}$$

Using Snell's law -
 $\sin i = \mu = \frac{\mu n / \mu B}{\mu B / \mu B}$

$$\mu n = \mu B \quad \text{put eq (1)}$$

$$A = \mu(\mu t \cos r + \mu B - \mu B)$$

$$\Delta = \mu(\mu t \cos r) \quad \text{--- (2)}$$

Now,

$$A = \mu(\mu t \cos r)$$

$$\Delta = \mu(\mu t \cos r) \quad \text{--- (3)}$$

in Δ $\mu B \cos r = \frac{\mu B}{\mu B}$

$$\mu B \cos r = \mu B \cos r$$

$$\mu B \cos r = (\mu t + \mu t) \cos r \quad \left\{ \begin{array}{l} \Delta AEO \cong \Delta BO \\ AO = BO = t \end{array} \right.$$

$$\mu B \cos r = \mu t \cos r \quad \text{put eq (3)}$$

$$\Delta = \mu t \cos r \quad \text{--- (4)}$$

(A) Interference in reflected system - The total path diff. in reflected system using stroke's treatment is given

$$\Delta T = \mu t \cos r - \frac{\mu}{2}$$

$$\Delta T = \mu t \cos r - \frac{\mu}{2}$$

(i) Condition for maxima - The intensity will be maximum -

$$\Delta r = n\lambda$$

$$2ut \cos r - d = n\lambda$$

$$2ut \cos r = (2n+1)\frac{d}{2}$$

(ii) Condition for minima - The intensity will be minimum when -

$$\Delta r = (2n+1)\frac{d}{2}$$

$$2ut \cos r - d = n\lambda - d$$

$$2ut \cos r = n\lambda$$

(B) Interference in transmitted system - The total path diff. in transmitted system

$$\Delta r = 2ut \cos r$$

(i) Condition for maxima - The intensity will be maxing -

$$2ut \cos r = n\lambda$$

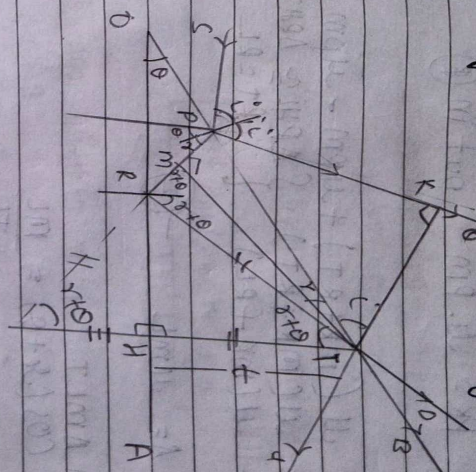
Condition for minima - The intensity will be minima -

$$\Delta r = \frac{(2n+1)\lambda}{2}$$

$$2ut \cos r = (2n+1)\frac{\lambda}{2}$$

Thus the condition for maxima or minima are complementary with each other in selected and transmitted system.

Interference due to wedge shape thin film -



When a thin film separating incident light is bounded by two inclined surfaces (AB) of varying thickness then it is called wedge shape thin film.

When a monochromatic light incident on the upper surface of wedge shaped film at point P then it suffers multiple reflection and refraction from the upper and lower surface of the film. The path difference between the rays is given by -

$$\Delta = \mu(PR + RT) - PR$$

$$\Delta = \mu(PR + RT) - PR$$

Using Snell's Law -

$$\frac{\sin i}{\sin r} = \mu = \frac{PK/PT}{PM/PT}$$

$$PK = \mu \cdot PM \quad \text{put in (3)}$$

$$\Delta = \mu [(mR + RT) + \mu PM - \mu PM]$$

$$\Delta = \mu [(mR + RT)] \quad \left\{ \begin{array}{l} \Delta RHL \approx \Delta RHT \\ RT = RL \end{array} \right.$$

$$\Delta = \mu mL \quad \text{--- (2)}$$

in ΔMLT

$$\cos(\beta + \theta) = \frac{mL}{TL}$$

$$mL = TL \cos(\beta + \theta)$$

$$mL = (TH + HL) \cos(\beta + \theta)$$

$$mL = (t + t) \cos(\beta + \theta) \quad \left\{ \begin{array}{l} \Delta RHL \approx \Delta RHT \\ TH = HL = t \end{array} \right.$$

$$mL = 2t \cos(\beta + \theta)$$

put in eq^y (2) we get

$$\Delta = 2\mu t \cos(\beta + \theta)$$

Performance in neglected system -

$$\Delta T = \frac{\Delta - d}{2}$$

$$\Delta T = 2\mu t \cos(\beta + \theta) - \frac{d}{2}$$

(1) Condition for maxima -

$$\Delta T = n\lambda \quad 2\mu t \cos(\beta + \theta) = \frac{d}{2} = n\lambda$$

$$2\mu t \cos(\beta + \theta) = (2n+1)\frac{\lambda}{2}$$

(2) Condition for minima -

$$\Delta T = n\lambda - \frac{d}{2}$$

$$2\mu t \cos(\beta + \theta) - \frac{d}{2} = n\lambda - \frac{d}{2}$$

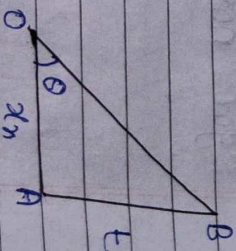
$$2\mu t \cos(\beta + \theta) = n\lambda$$

Fringe width -

Δx_{HB}

$$\tan \theta = \frac{p}{B} = \frac{t}{x_n}$$

$$t = x_n \tan \theta$$



The condition for minima neglected system

$$2\mu t \cos(\beta + \theta) = n\lambda$$

$$\text{then } 2\mu x_n \tan \theta \cos(\beta + \theta) = n\lambda \quad \text{--- (1)}$$

The condition for (n-1)th minima in neglected system -

$$2\mu t x_{n-1} \tan \theta \cos(\beta + \theta) = (n-1)\lambda \quad \text{--- (2)}$$

Now, (1) - (2) we have

$$2\mu x_n \tan \theta \cos(\beta + \theta) [x_n - x_{n-1}] = n\lambda - (n-1)\lambda$$

The circular concentric interference fringes are observed when a thin film of air or other transparent medium is enclosed b/w plane glass plate and plane convex lens such fringes are known as Newton's ring

When a monochromatic light incident normally on the upper surface of plane convex lens. Then the ray reflected from upper and lower surface of film in such a way that it produce circular rings.

Newton's ring a circular because the focus of constant thickness about the point of contact circular. The centre of Newton's rings in reflected system is dark because path diff. in reflected system

$$\Delta = 2\mu t - \frac{d}{2}$$

$$t = 0$$

$$\Delta = -\frac{d}{2}$$

which is the condition of minima.

So, centre becomes dark.

Let CME be vertical section of plane convex lens.

$$CD \times DE = AD \times DM$$

$$CD \times DE = (R - DM) \times DM$$

Let d_n and d_n be the radii and diameter of n th ring pass through C and E.

$$d_n \times d_n = (2R - t) \times t$$

$$d_n^2 = 2Rt - t^2$$

$$d_n^2 = \frac{2Rt}{R}$$

$$d_n = \frac{\sqrt{2Rt}}{R}$$

$$d_n = \frac{d_n^2 \sqrt{R}}{2R}$$

$$\left\{ \begin{array}{l} t \ll R \\ \text{then } t^2 \approx 0 \end{array} \right\}$$

$$\left\{ \begin{array}{l} d_n = \frac{d_n^2}{2} \\ \end{array} \right\}$$

The path difference b/w interfering ray is given by - when $\theta \ll \epsilon$ for normal incidence

(A) Interference in reflected system - The total path difference in reflected system is given by using Stokes treatment in given

$$\Delta r = 2\mu t - \frac{d}{2}$$

for maxima -

$$\Delta r = n\lambda$$

$$2\mu t - \frac{d}{2} = n\lambda$$

$$2\mu t = (n + \frac{1}{2})\lambda$$

for air medium $\mu = 1$

$$\Delta t = (n + \frac{1}{2})\lambda$$

$$\frac{d^2n}{4R} = (n + \frac{1}{2})\lambda$$

$$d^2n = 4R\lambda (n + \frac{1}{2})$$

$$dn \propto \sqrt{(n + \frac{1}{2})} \text{ Bright}$$

Thus the diameter of n^{th} bright ring in reflected system is directly proportional to square root of half of odd five integer.

(ii) for minimum intensity will be minimum

$$\Delta t = \frac{d}{2} = n\lambda$$

$$\Delta t = n\lambda$$

for air $w=1$

$$\Delta t = n\lambda$$

$$\frac{d^2n}{4R} = n\lambda$$

$$d^2n = 4R\lambda n$$

$$dn \propto \sqrt{n}$$

Dark

Thus the diameter of n^{th} dark ring in reflected system is directly proportional to square root of positive integer.

(B) Interference in transmitted system - The path diff. in transmitted system is given by

$$\Delta T = \Delta t$$

(i) Condition for maxima -

$$\Delta T = n\lambda$$

$$\Delta t = n\lambda$$

for air $w=1$

$$\Delta t = n\lambda$$

$$\frac{d^2n}{4R} = n\lambda$$

$$d^2n = 4R\lambda n$$

Bright

(ii)

for minimum

$$\Delta t = (n + \frac{1}{2})\lambda$$

$$\Delta t = (n + \frac{1}{2})\lambda$$

for air $w=1$

$$\Delta t = (n + \frac{1}{2})\lambda$$

$$\frac{d^2n}{4R} = (n + \frac{1}{2})\lambda$$

$$dn \propto \sqrt{n + \frac{1}{2}}$$

Dark

Determination of wavelength using Newton's Ring - The reflected system for air medium is given by -

$$d_n^2 = 4Rdn \quad \text{--- (1)}$$

Condition for n^{th} minima in reflected system for air medium is given by

$$d_n^2 = 4mRd \quad \text{--- (2)}$$

Now, (1) - (2)

$$d_n^2 - d_n^2 = 4Rd(m-n)$$

Let

$$\boxed{m-n = p}$$

$$d_{m+p}^2 - d_n^2 = 4Rd \cdot 4pRd$$

$$\boxed{d = \frac{d_{m+p}^2 - d_n^2}{4pR}}$$

determination of refractive index - The condition for n^{th} minima in reflected system in a medium of refractive index μ is given by -

$$\mu d_n^2 = 4nRd \quad \text{--- (1)}$$

The condition for n^{th} minima in reflected system for air medium is given by

$$d_n^2 \text{air} = 4nRd \quad \text{--- (2)}$$

eq (1) / 2

$$\frac{\mu d_n^2}{d_n^2 \text{air}} = 1$$

$$\boxed{\mu = \frac{d_n^2 \text{air}}{d_n^2 \text{refl}}}$$

Q. Calculate the thickness of the thin film ($\mu = 1.4$) in which interference of violet component ($\lambda = 4000 \text{ \AA}$) of incident light can take place by reflection?

Ans.

$$2t \cos r = (2n+1) \frac{\lambda}{2}$$

$$t = \frac{(2n+1) \lambda}{4 \cos r}$$

$$t_{\text{min}} = \frac{\lambda}{4 \mu}$$

$$t_{\text{min}} = \frac{4000}{4 \times 1.4} = 714.28 \text{ \AA}$$

$$\boxed{t_{\text{min}} = 714.28 \text{ \AA}}$$

Q. A man whose eyes 150 cm above the ground level observe the greenish colour at a distance of 100 cm from his feet. Calculate the possible thickness of film.

Ans:

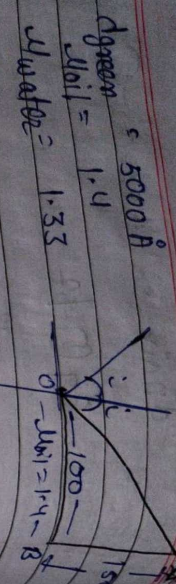
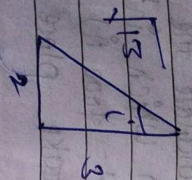


Diagram $\mu = 5000 \text{ \AA}$
 $\mu_{\text{air}} = 1.4$
 $\mu_{\text{water}} = 1.33$

2 ut $\cos r = (2m+1) \frac{\lambda}{2}$
 $t = \frac{(2m+1) \lambda}{2 \mu \cos r}$

for $m = 1$ $\tan i = \frac{100}{150} = \frac{2}{3}$

$\sin i = \frac{2}{\sqrt{13}}$



$\sin i' = \mu$

$\sin r = \frac{\sin i'}{\mu}$

$\cos r = \sqrt{1 - \frac{\sin^2 i'}{\mu^2}} = \sqrt{1 - \frac{4}{13}}$
 $\cos r = 0.9182$

$t = \frac{5000 \times 10^{-8}}{2 \times 1.4 \times 0.9182} (2m+1)$

$t = 9.7125 \times 10^6 (2m+1) \text{ cm}$ Ans.
 $t = 9.725 \times 10^6 \text{ cm}$

Q. White light is incident on a soap film at an angle $\sin^{-1} \frac{4}{5}$ and the reflected light is observed with spectroscope. It is found with two consecutive dark bands $6 \times 10^5 \text{ cm}$. If the refractive index of the film is $\frac{4}{3}$ calculate the thickness?

Ans:

2 ut $\cos r = n \lambda_1$ --- (1)
 $2 ut \cos r = (n+1) \lambda_2$ --- (2)

$(n+1) \lambda_2 = n \lambda_1$
 $\lambda_2 = n \lambda_1 - n \lambda_2$
 $n = \frac{\lambda_2}{\lambda_1 - \lambda_2}$

2 ut $\cos r = \frac{\lambda_1 - \lambda_2}{\lambda_1 - \lambda_2}$

$t = \frac{\lambda_1 \lambda_2}{2 \mu \cos r (\lambda_1 - \lambda_2)}$

$\cos r = \sqrt{1 - \left(\frac{\sin^2 i'}{\mu^2} \right)}$

$\cos r = \sqrt{1 - \frac{(4/5)^2}{(4/3)^2}} = \frac{4}{5}$

$t = \frac{6.1 \times 10^5 \times 6.0 \times 10^5}{2 \times \frac{4}{3} \times \frac{4}{5} \times 0.1 \times 10^5} = 0.0017 \text{ cm}$

Q. Light of wavelength 600 Å falls normally on a thin wedge of glass forming fringes of equal thickness. Find the angle of wedge in seconds.

Solⁿ.

$\mu = 1.4$
 $\beta = 2.0 \text{ mm}$

$\beta = \frac{\Delta}{2\mu\theta}$

$\theta = \frac{\Delta}{2\mu\beta}$

$\theta = \frac{1}{2 \times 1.4} \times \frac{180}{\pi} \times 60 \times 10^{-6} \text{ sec}$

$\theta = \frac{6000 \times 10^{-6}}{2 \times 1.4 \times 2.0 \times 10^{-3}} \times \frac{180}{\pi} \times 60 \times 10^{-6}$

$\theta = 22.09 \text{ sec}$

Q. Newton's Ring are observed normally in reflected light of wavelength 6000 Å the diameter of a 10th dark ring is 0.50 cm find the radius of curvature of the lens and the thickness of film.

$\lambda = 6000 \text{ Å}$
 $d_{10} = 0.50 \text{ cm}$

$2\mu t = n\lambda$

for $\mu = 1$

$2t = n\lambda$
 $\frac{d_n^2}{4R} = n\lambda$

$R = \frac{d_n^2}{4n\lambda}$
 $R = \frac{(0.50)^2}{4 \times 10 \times 6000 \times 10^{-8}}$
 $R = 104 \text{ cm}$

(ii)
 $2t = \frac{d_n^2}{4R}$

$t = \frac{d_n^2}{8R} = \frac{(0.50)^2}{8 \times 104}$

$t = 5 \times 10^{-4} \text{ cm}$

Q. Newton's ring are observed by keeping a spherical surface of 100 cm radius on a plane glass plate the diameter of 15th bright ring is 0.89 cm and the diameter of 5th ring is 0.336 cm what is the wavelength of light used.

$\lambda = \frac{d_n^2 - d_p^2}{4PR}$

$R = 100 \text{ cm}$

$d_{15} = 0.89 \text{ cm}$

$d_5 = 0.336 \text{ cm}$

$d = ?$

$\lambda = \frac{d_{15}^2 - d_5^2}{4 \times P \times R} = \frac{(0.89)^2 - (0.336)^2}{4 \times P \times 100}$

$\left. \begin{matrix} P = 15 \\ P = 5 \end{matrix} \right\}$

$d = 5880 \text{ Å}$

Q. In Newton's ring experiment the diameter of the 4th & 12th dark ring are 0.400 cm & 0.700 cm respectively deduce the diameter of the 8th dark ring

Sol.

$$d_1 = 0.400 \text{ cm}$$

$$d_2 = 0.700 \text{ cm}$$

$$d_{20} = ?$$

$$\text{Ans. } \lambda = \frac{d_{n+1}^2 - d_n^2}{4nR}$$

$$d_{n+1}^2 - d_n^2 = 4nR\lambda \quad \text{--- (1)}$$

$$d_{12}^2 - d_4^2 = 4 \times 8 \times R \lambda \quad \text{--- (2)}$$

$$d_{20}^2 - d_{12}^2 = 4 \times 8 \times R \lambda \quad \text{--- (3)}$$

$$\text{Eq (3)} \div \text{Eq (2)} \Rightarrow \frac{d_{20}^2 - d_{12}^2}{d_{12}^2 - d_4^2} = 1$$

$$d_{20}^2 - d_{12}^2 = d_{12}^2 - d_4^2$$

$$d_{20}^2 = 2(d_{12}^2) - (d_4^2)$$

$$d_{20} = 0.82$$

$$\boxed{d_{20} = 0.90}$$